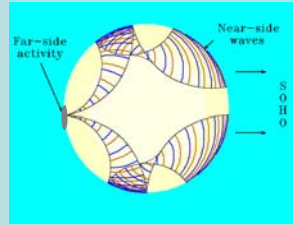
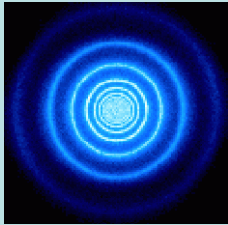


## ASTR 7500: Solar & Stellar Magnetism

Hale CGEP Solar & Space Physics



Profs. Brad Hindman & Juri Toomre

Lecture 24 Thurs 18 Apr 2013

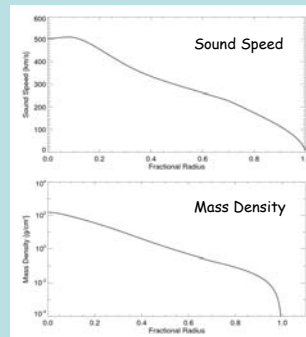
[zeus.colorado.edu/astr7500-toomre](http://zeus.colorado.edu/astr7500-toomre)

## Lecture 24 Helioseismic Discoveries

- Results from Global Helioseismology
  - Sound speed and density inversions
  - Rotation rate inversion
- Local Helioseismology
  - Ring analysis
  - Meridional Flows
  - Flows around active regions and sunspots
  - Farside imaging

## Results from Global Helioseismology

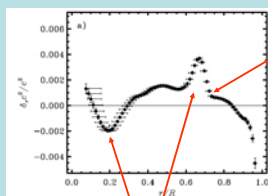
## Internal Thermodynamics



Solar model and an inversion from observations

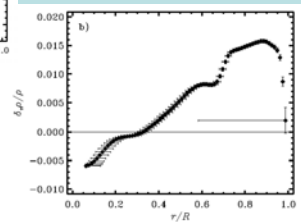
Since the solar model used in the inversion is very good (predicts frequencies to within a percent or so of the observed), the inversion results and the model are indistinguishable in plots of this nature.

## Closer Look



Bottom of the convection zone

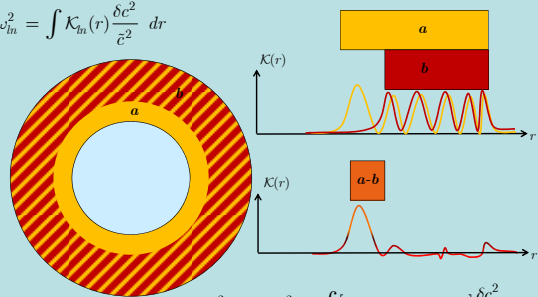
Error  
Depth Resolution



Likely due to weak mixing of elements (or perhaps errors in the model opacity).

## Overlapping Kernels

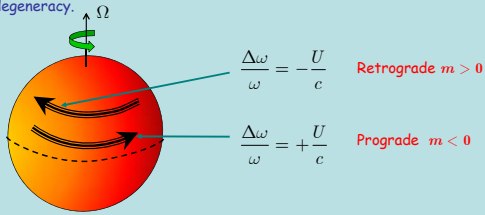
$$\delta\omega_{in}^2 = \int \mathcal{K}_a(r) \frac{\delta c^2}{c^2} dr$$



$$\delta\omega_a^2 - A_b \delta\omega_b^2 = \int [\mathcal{K}_a(r) - A_b \mathcal{K}_b(r)] \frac{\delta c^2}{c^2} dr$$

## Solar Rotation

Remember that the eigenfrequencies were degenerate in the quantum number  $m$ , for a spherically symmetric star. Rotation breaks that degeneracy.



$$\frac{\Delta\omega}{\omega} = -\frac{U}{c} \quad \text{Retrograde } m > 0$$

$$\frac{\Delta\omega}{\omega} = +\frac{U}{c} \quad \text{Prograde } m < 0$$

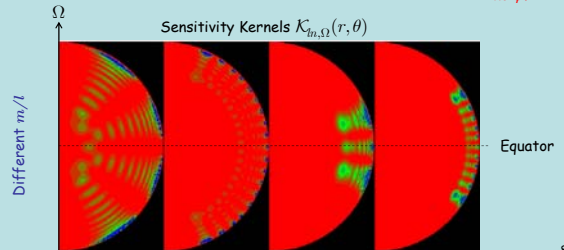
$$W(r) Y_l^m(\theta, \varphi) e^{-i\omega t} \sim W(r) P_l^m(\cos\theta) e^{i(m\varphi - \omega t)}$$

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## Differential Rotation

Rotation causes a splitting of the frequency that varies with  $m$ . As before, the splitting can be expressed using a kernel

$$\omega_{nlm} - \omega_{nl0} = \Delta\omega_{nlm} = m \iint \mathcal{K}_{m,\Omega}(r, \theta) \Omega(r, \theta) r dr d\theta$$



8

## Inversion for Rotation

$$\omega_{nlm} - \omega_{nl0} = \Delta\omega_{nlm} = m \iint \mathcal{K}_{m,\Omega}(r, \theta) \Omega(r, \theta) r dr d\theta$$

An inversion for the rotation rate proceeds just as for the sound speed. We parameterize the rotation rate with a set of free parameters  $\alpha_j$ .

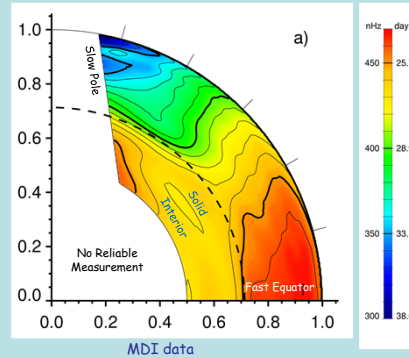
$$\Omega(r, \theta) \approx \Omega_{\text{mod}}(r, \theta; \alpha)$$

We find the parameters  $\alpha_j$  that best fit the data by minimizing . . .

$$\chi^2(\alpha) = \sum_{nlm} \frac{\left[ \Delta\omega_{nlm} - m \iint \mathcal{K}_{m,\Omega}(r, \theta) \Omega_{\text{mod}}(r, \theta; \alpha) r dr d\theta \right]^2}{\sigma_{nlm}^2} + \lambda \iint |\nabla^2 \Omega_{\text{mod}}(r, \theta; \alpha)|^2 r dr d\theta$$

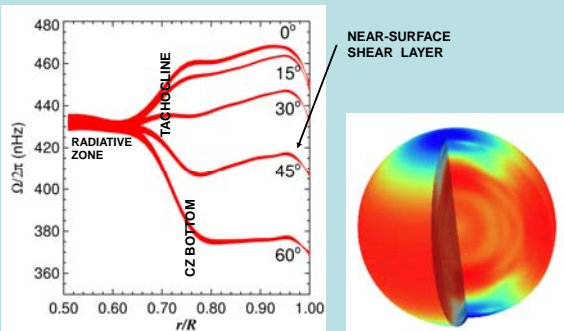
9

## Differential Rotation in the Sun



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## Shear Layers



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## Meridional Circulation

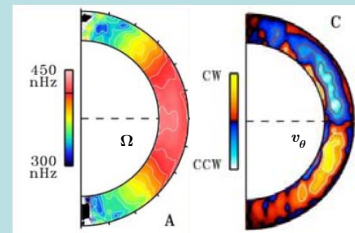
### Flux Transport Dynamos

- Carries the poloidal flux from the surface to the tachocline.
- Sets the period of the activity cycle.

Miesch et al. 2008

### Numerical Simulations

- Helioseismic constraints on the differential rotation have been invaluable guides.
- Meridional circulations could do the same.



## Bah!

Unfortunately, the *global modes* are insensitive to a meridional flow.

### Why?

- The global modes are standing waves in latitude (unlike in longitude)
- They spend an equal amount of time swimming upstream and downstream.

$$W(r) P_l^m(\cos \theta) e^{i(m\phi - \omega t)}$$

↑ Real function (standing wave)     ↑ Complex (propagating)

The global modes also have the weakness that their energy densities (or the sensitivity kernels) are:

- Symmetric about the equator (insensitive to hemispheric differences)
- Constant with longitude (insensitive to longitudinal variation)

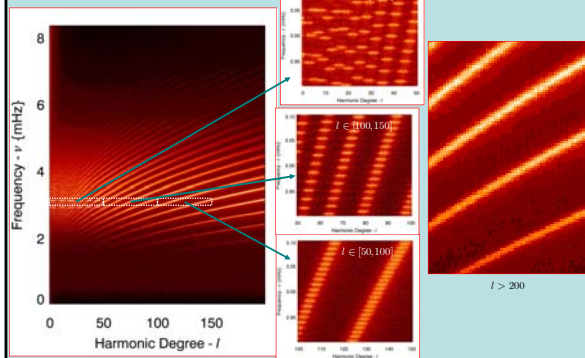
$$\left[ W(r) P_l^m(\cos \theta) \right]^2 \left| e^{i(m\phi - \omega t)} \right|^2$$

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## Local Helioseismology

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## Peaks to Ridges



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## Wavenumber Regimes

### Low-degree waves ( $l < 200$ )

- Weakly damped and able to circumnavigate the sun.
- Form resonances → integer  $l$
- Isolated peaks
- The frequencies are global measures of the internal properties
- Used by Global Helioseismology

### High-degree waves ( $l > 200$ )

- Highly damped and unable to travel completely around the sun.
- All horizontal wavenumbers are allowed
- Ridges of merged peaks
- The frequencies are local measures of the internal properties
- Used by Local Helioseismology

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## Two Primary Flavors

Local helioseismology utilizes the high-degree waves to make local estimates of solar properties. There are two primary flavors (some with nuts, others with chocolate chips, etc.)

### Correlation Procedures

(*Time-Distance Tomography, Acoustic Holography*)

- The fundamental measurement is the time it takes a signal to travel from one point on the surface to a second point.
- Such measurements are made by cross-correlating the two time series associated with those points and fitting for the time lag associated with maximum correlation.

### Spectral Procedures

(*Ring Analysis, Fourier-Hankel Decomposition*)

- The fundamental measurement is the frequencies of local modes.
- Such measurements are made by fitting peaks in a power spectrum.

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## Ring Analysis

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## Plane Wave Decomposition

Each image is apodized and then transformed

$$v(k_x, k_y, \omega) = \iiint e^{-i\mathbf{k}_s \cdot \mathbf{r}} e^{i\omega t} v_{\text{obs}}(x, y, t)$$

**Why plane waves?**  
Waves of short wavelength (larger wavenumber) are confined near the surface.

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## Ring Spectra

Power =  $|v(k_x, k_y, \omega)|^2$

Since there are two spatial dimensions ( $x, y$ ) and the time dimension  $t$ , the spectra are 3D.

Global Spectra

Local Spectra

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## Rings!

The modes appear as nested trumpets aligned with the frequency axis.

Cuts at constant frequency produce nested rings that have larger radii for larger frequencies.

Changes in the mode frequencies manifest as changes to the radii of the rings. Therefore, careful measurement of the radii can be used to determine subsurface structure ( $c^2, \rho$ , etc.).

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## The effects of a wind

In a 2-D homogeneous fluid with sound speed  $c$ , the dispersion relation for acoustic waves is,

$$\omega^2 = k^2 c^2 = (k_x^2 + k_y^2) c^2$$

No Wind

Circular cross-section at constant frequency

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## Doppler Shift Caused by a Wind

A wind with speed  $U$  in the  $x$  direction causes a Doppler shift. The shift can be expressed using a Galilean transformation

$$(\omega - k_x U)^2 = k^2 c^2$$

With Wind

Elliptical cross-section at constant frequency

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## Effect on $p$ -Mode Spectra

**Tracked**

The above spectra was obtained by following the same patch of fluid as it rotates across the solar disk. This removes the large rotational velocity.

**Not Tracked**

The above spectra was obtained by studying the same area on the solar disk. Equatorial rotation results in a speed of  $\sim 2000$  m/s.

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## Measuring Frequencies and Doppler Shifts

The frequency of a mode and its Doppler shift are determined by fitting the spectra with a "line profile", usually a Lorentzian function of frequency. This fitting is performed for each value of  $k_h$  in the spectra.

$$P(\vec{k}_h, \omega) = \frac{A}{(\omega - \omega_n + \vec{k}_h \cdot \vec{U})^2 + \Gamma^2}$$

$$\vec{U} = U_x \hat{x} + U_y \hat{y}$$

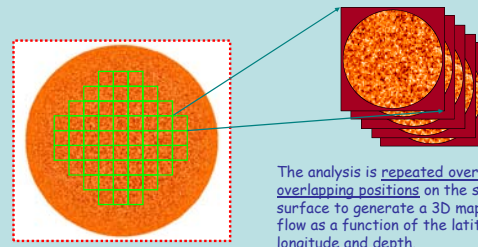
$A$  - Amplitude  
 $\Gamma$  - Line widths  
 $U_x$  and  $U_y$  - Doppler shifts  
 $\omega_n$  - Mode frequency

The frequencies  $\omega_n$  are a measure of the sound speed below the surface.

The Doppler shifts  $U_x$  and  $U_y$  measure the flow speed in the  $x$  and  $y$  directions.

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## Building up a map



The analysis is repeated over many overlapping positions on the solar surface to generate a 3D map of the flow as a function of the latitude, longitude and depth.

The analysis is repeated on a daily basis to build up a time series of 3D flows.

The size of the analysis tile determines the horizontal resolution

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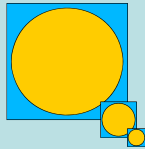
## 3-D Inversion

- RLS (Regularized Least Squares)

$$\chi^2 = \sum_i \frac{1}{\sigma_i^2} [U_{x,i} - \iiint d^3x K_i(\vec{x}) v_x(\vec{x})]^2 + \lambda \iiint d^3x |\vec{\nabla} v_x|^2$$

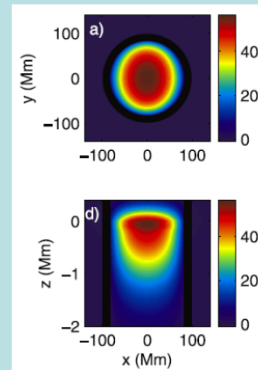
- Employs Three Tile Sizes

- $2^\circ \times 2^\circ$  (22 Mm in diameter)
- $4^\circ \times 4^\circ$  (45 Mm in diameter)
- $16^\circ \times 16^\circ$  (183 Mm in diameter)



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## Sensitivity Kernels



$f$  mode

Horizontal variation of a sensitivity kernel.

The circular shape is determined by the apodization.

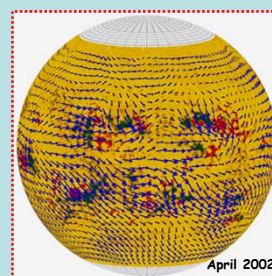
Vertical variation of a sensitivity kernel.

The depth dependence is determined by eigenfunctions energy density.

## Flows Measured with Local Helioseismology

## Large-Scale Flow Maps

### Synoptic Flow Map



7 Mm below the surface.

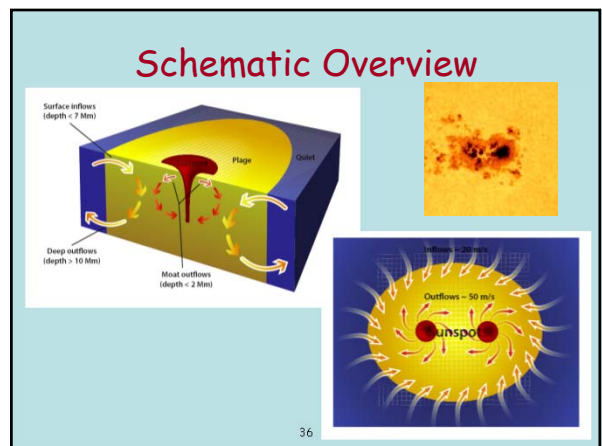
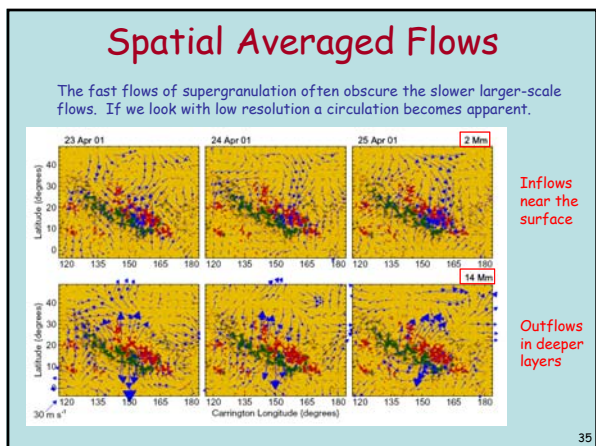
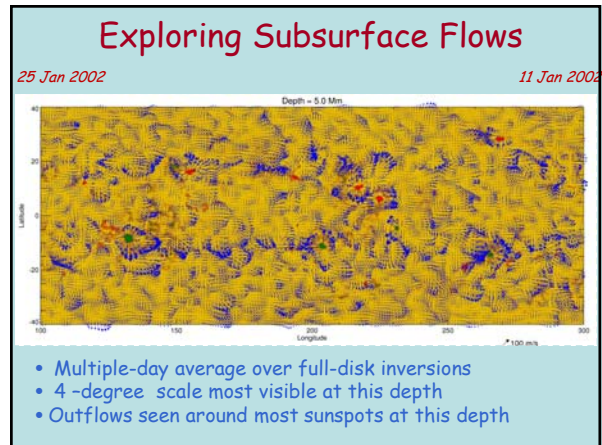
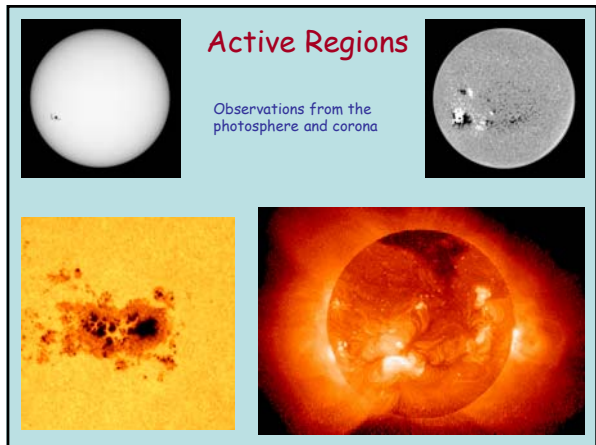
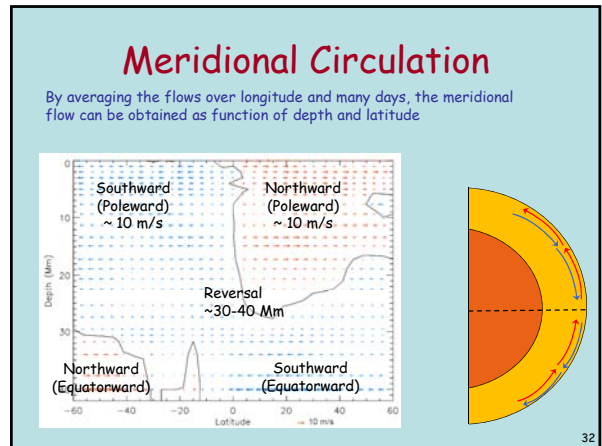
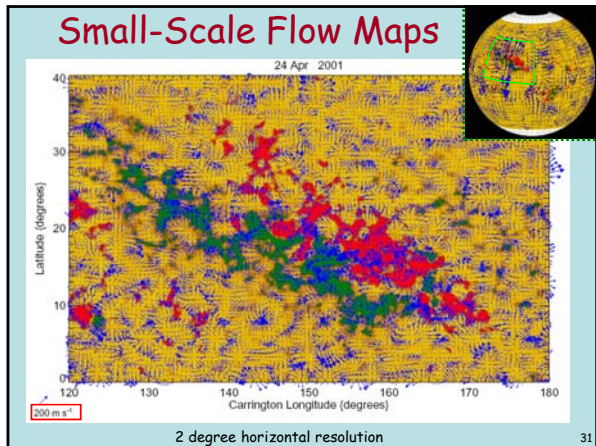
Average flow over a month.

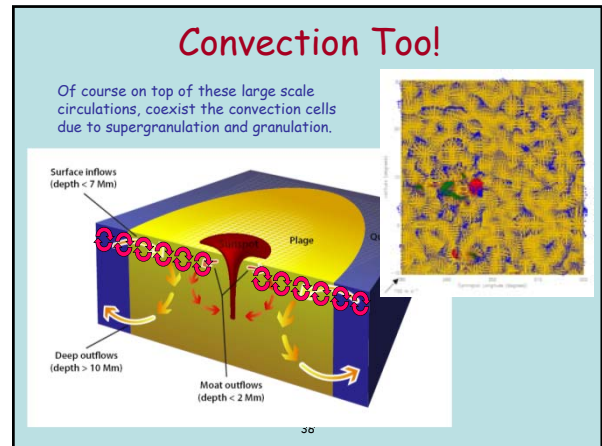
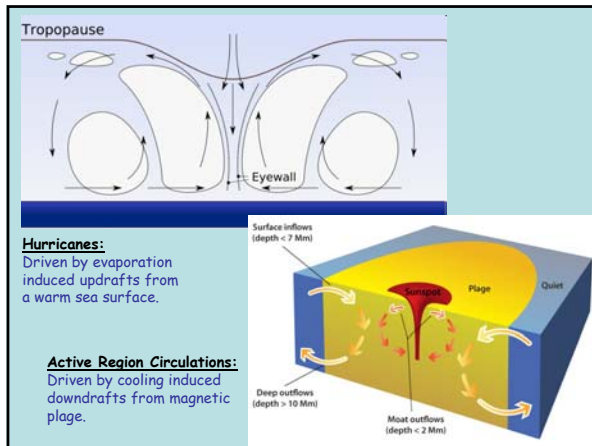
The surface differential rotation has been removed.

Flow has been averaged spatially over a scale of  $\sim 150$  Mm.

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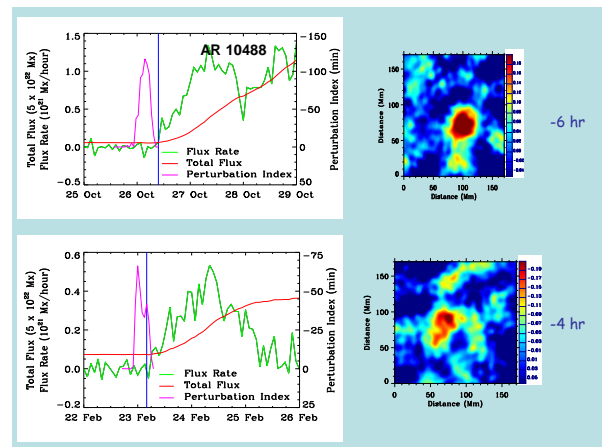
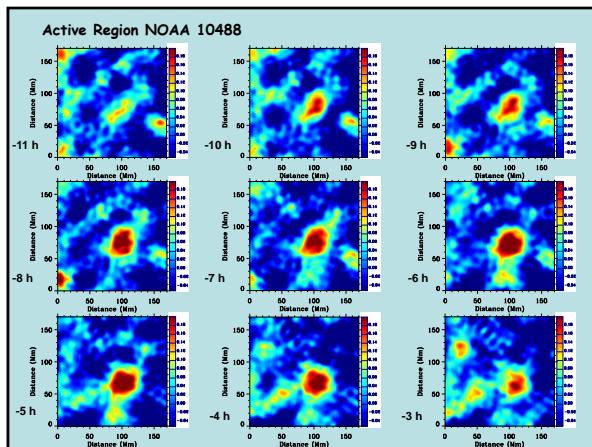
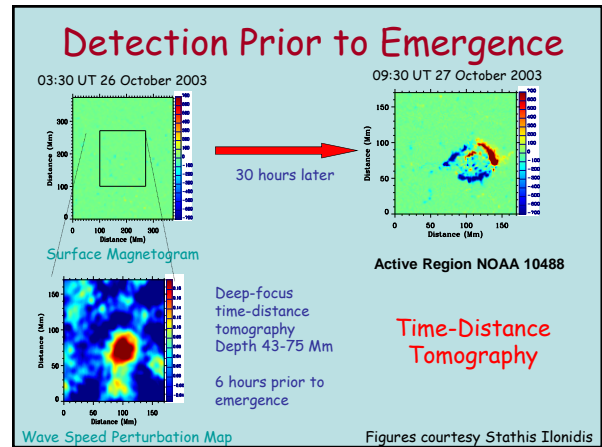




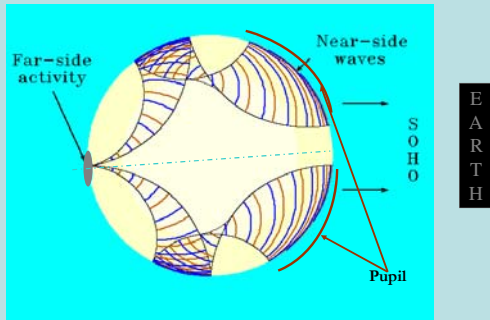


## Detecting Active Regions (Before they can be seen at the surface)

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## Farside Imaging



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## Farside Maps

